

Reminder from the previous lecture:

**d-regular graphs :**

Graphs where all the vertices have the same degree.

Examples:

1. Complete graphs
2. Lattices
3. d-regular infinite trees.

**d-regular graphs are an expanding family.**

$$E = \frac{Vd}{2} \quad \rightarrow \quad V \text{ or } d \text{ must be even}$$

$$\beta = V\left(\frac{d}{2} - 1\right)$$

$\mathcal{G}_{V,d}$ . The ensemble of all the  $d$  regular graphs with  $V$  vertices .

$$\text{For fixed } d \text{ and } V \rightarrow \infty : |\mathcal{G}_{V,d}| \approx \sqrt{2} e^{\frac{1-d^2}{4}} \left( \frac{d^d V^d}{e^d (d!)^2} \right)^{\frac{V}{2}}$$

$\langle \dots \rangle_{\mathcal{G}}$  : Ensemble average taken with uniform probability distribution.

$C_t$  : Number of  $t$  - periodic cycles with no back tracking

$$\langle C_t \rangle_{\mathcal{G}} = \frac{V \cdot d \cdot (d-1)^{t-1} \cdot \frac{d-1}{dV}}{2t} = \frac{(d-1)^t}{2t}$$

Hence, short  $t$  - periodic cycles with  $t < \log_{d-1} V$  are **rare** .

The  $R$ -neighbourhood of every vertex for  $R \leq \log_{d-1} \frac{V}{2}$  is almost surely a  $d$ -regular tree.

The diameter of a  $G(V, d)$  graph, i.e. the maximal distance between vertices in  $G$ , is given by

$$\text{diam}(G) = \log_{d-1}(V \log_{d-1} V) + O(1)$$

Therefore, the typical distance between vertices along the boundary of the 'local tree' is of the same magnitude as the distance between two arbitrary vertices in  $G$ .

Denote by  $C_t$  the number of  $t$ -cycles (primitive, non backscatter, non self intersecting  $t$ -periodic orbits).

For  $t < \log_{d-1} \frac{V}{2}$ , the  $C_t$  distribute as independent Poisson variables with a mean:  $\langle C_t \rangle_G = \frac{(d-1)^t}{2t}$

**Reminder :** For  $d$ -regular graphs the spectra of  $A$  and  $L$  are the same but for a shift by  $d$  and sign change.

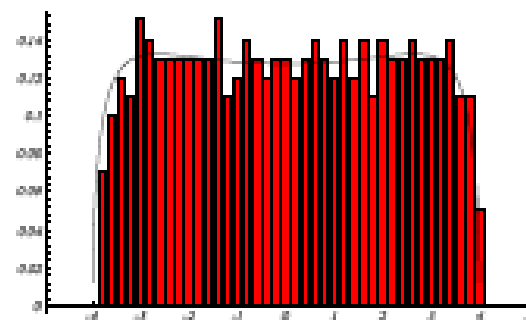
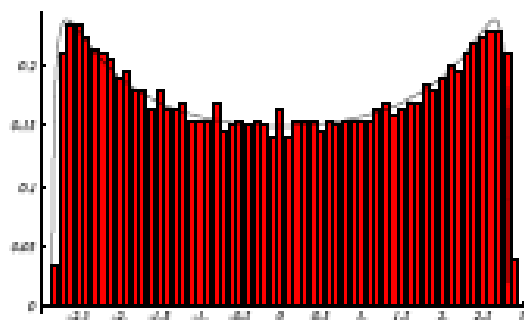
## Spectral properties of d-regular graphs

The spectral density  $\rho(\lambda) = \frac{1}{V-1} \sum_{k=1}^{V-1} \delta(\lambda - \lambda_k)$  (The trivial eigenvalue is excluded)

**The mean spectral density for  $d$  regular graphs, ( $V \rightarrow \infty, d = \text{const}$ )**

Kesten MacKay limit distribution: (Supported in  $|\lambda| < 2\sqrt{d-1}$ ):

$$\rho_{KM}(\lambda) = \lim_{V \rightarrow \infty} \frac{1}{V-1} \left\langle \sum_{k=1}^{V-1} \delta(\lambda - \lambda_k) \right\rangle_{\mathcal{G}} = \frac{d}{2\pi} \frac{\sqrt{4(d-1) - \lambda^2}}{d^2 - \lambda^2}$$



(a) Cubic graph on 2000 vertices. (b) 5-valent graph on 500 vertices.

Figure 1. Eigenvalue distributions of random graphs vs McKay's law

# EIGENVALUE SPACINGS FOR REGULAR GRAPHS

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Unfolding the spectrum with the Kesten-McKay density using  $\mathcal{N}_{KM}$  : The mean spectral counting function.

$$s_j = \mathcal{N}_{KM}(\lambda_j) \quad ; \quad \frac{d\mathcal{N}_{KM}}{d\lambda} = \rho_{KM}(\lambda) \quad ; \quad ds = \rho_{KM}(\lambda)d\lambda = \frac{d\lambda}{\langle d\lambda \rangle}$$

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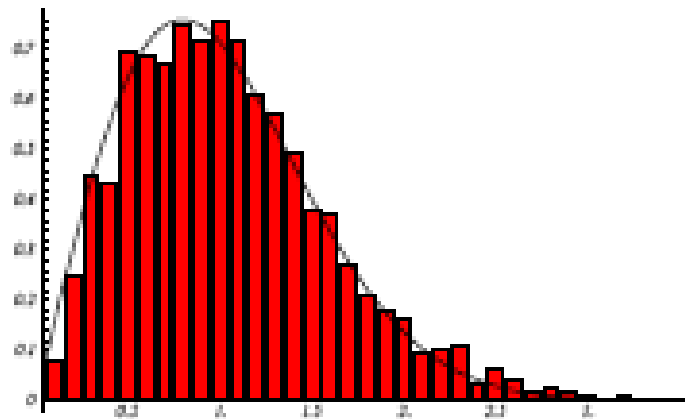


Figure 2. Level spacing distribution of a cubic graph on 2000 vertices vs GOE

## 2. The trace formula

Purpose: express the spectral density in terms of "periodic orbits"

Connectivity of directed edges

$$B_{e,e'} = \delta_{o(e),\tau(e')}$$

$$J_{e,e'} = \delta_{o(e),\tau(e')} \delta_{o(e'),\tau(e)}$$

$$Y = B - J$$

$Y$  excludes back tracking connections!

Note: The number of non-backtracking periodic walks =  $\text{tr}Y^t$

However: The number of all periodic walks =  $\text{tr}A^t = \text{tr}B^t$

## Bartholdi's Identity for $d$ -regular graphs

Define the two  $2E \times V$  rectangular matrices

$$B_{e,i}^{(+)} := \begin{cases} 1 & \text{if } t(e) = i \\ 0 & \text{otherwise} . \end{cases} \quad ; \quad B_{e,i}^{(-)} := \begin{cases} 1 & \text{if } o(e) = i \\ 0 & \text{otherwise} . \end{cases}$$

Denoting by  $\widetilde{X}$  the transpose of  $X$ , one can easily prove that

$$\begin{aligned} B^{(+)} \widetilde{B^{(-)}} &= B & ; & \quad \widetilde{B^{(-)}} B^{(+)} = A \\ \widetilde{B^{(+)}} B^{(+)} &= dI^{(V)} & ; & \quad B^{(+)} \widetilde{B^{(+)}} = YJ + I^{(2E)} . \end{aligned} \quad (1)$$

For arbitrary complex  $s, w$  construct the two  $(2E+V) \times (2E+V)$  square matrices

$$L = \begin{bmatrix} (1 - w^2 s^2) I^{(V)} & -\widetilde{B^{(-)}} + w s \widetilde{B^{(+)}} \\ 0 & I^{(2E)} \end{bmatrix} ; \quad M = \begin{bmatrix} I^{(V)} & \widetilde{B^{(-)}} - w s \widetilde{B^{(+)}} \\ s B^{(+)} & (1 - w^2 s^2) I^{(2E)} \end{bmatrix}$$

Using the identities (1) one can compute the matrices  $LM$  and  $ML$ , and since their determinants are equal, one finally gets Bartholdi's identity.

$$\det(I^{(2E)} - s(B - wJ)) = (1 - w^2 s^2)^{E-V} \det(I^{(V)}(1 + w(d - w)s^2) - sA) .$$

## 2.a The “canonical “ trace formula ( $w = 1$ )

Bass’ (Bartholdi’s) identity :

$$\det(\eta I^{(2E)} - Y) = (\eta^2 - 1)^{E-V} \det(I^{(V)}(\eta^2 + (d - 1)) - \eta A) .$$

$\eta$  arbitrary complex number.

$I^{(2E)}$  and  $I^{(V)}$  unit matrix in  $2E$  and  $V$  dimensions.

$A_{i,j}$  : Adjacency Matrix,  $\dim A = V$  ;  $B_{e',e} = \delta_{o(e'),\tau(e)}$  ,  $\dim B = 2E$ .

$J_{e',e} = \delta_{e',\hat{e}}$  ;  $Y = B - J$  : The Hashimoto connectivity matrix.

$$\#\{t\text{-periodic walks}\} = \text{tr} B^t = \text{tr} A^t.$$

$$\#\{t\text{-periodic non-backscattering walks}\} = \text{tr} Y^t$$

**The spectrum of A can be used to count n.b. periodic walks  
and vice versa**

**The counts of n.b. periodic walks can be used to compute the spectrum of A**

Bartholdi again :  $\det(\eta I^{(2E)} - Y) = (\eta^2 - 1)^{E-V} \det(I^{(V)}(\eta^2 + (d-1)) - \eta A)$  .

The spectrum of  $Y$ :  $\sigma(Y) =$

$$\left\{ 1 \times (E - V), (-1) \times (E - V), \{\eta_l^\pm : \eta_l^{\pm 2} + (d-1) - \lambda_l \eta_l^\pm = 0, l = 1, \dots, V\} \right\}$$

$$\eta_l^\pm = \frac{1}{2} \left( \lambda_l \pm i \sqrt{4(d-1) - \lambda_l^2} \right) .$$

$$\lambda_0 = d \quad \rightarrow \quad \eta_0^\pm = (d-1), 1$$

**Assume Ramanujan:**  $\eta_l^\pm = \sqrt{d-1} e^{\pm i\phi_l}$  ;  $\phi_l = \arccos \frac{\lambda_l}{2\sqrt{d-1}}$   
 $l = 1, \dots, V-1$

**The number of t-periodic non backscattering walks :**

$$\text{tr } Y^t = (d-1)^t + (E-V)(1 + (-1)^t) + 1 + 2(d-1)^{\frac{t}{2}} \sum_{l=1}^{V-1} \cos t\phi_l .$$

Remember  $E = \frac{Vd}{2}$  and use Chebyshev Polynomials :  $T_n(x) = \cos(n \arccos(x))$

$$\text{tr } Y^t = (d-1)^t + V\left(\frac{d}{2} - 1\right)(1 + (-1)^t) + 1 + 2(d-1)^{\frac{t}{2}} \sum_{l=1}^{V-1} T_t \left( \frac{\lambda_l}{2\sqrt{d-1}} \right) .$$



$$y_t = \frac{1}{V} \frac{\text{tr } Y^t - (d-1)^t - 1}{(d-1)^{\frac{t}{2}}} = \frac{\left(\frac{d}{2} - 1\right)}{(d-1)^{\frac{t}{2}}} (1 + (-1)^t) + \frac{2}{V} \sum_{l=1}^{V-1} T_t \left( \frac{\lambda_l}{2\sqrt{d-1}} \right).$$

Multiply by  $\frac{1}{\pi(1+\delta_{t,0})} T_t \left( \frac{\lambda}{2\sqrt{d-1}} \right)$ , and sum over  $t$ :

$$\begin{aligned} & \frac{1}{\pi} \sum_{t=0}^{\infty} \frac{1}{1 + \delta_{t,0}} T_t \left( \frac{\lambda}{2\sqrt{d-1}} \right) y_t = \\ & = \frac{d-2}{2\pi} \sum_{t=0}^{\infty} \frac{1}{1 + \delta_{t,0}} \frac{(1 + (-1)^t)}{(d-1)^{\frac{t}{2}}} T_t \left( \frac{\lambda}{2\sqrt{d-1}} \right) + \frac{1}{V} \sum_{l=1}^{V-1} \delta_T \left( \frac{\lambda - \lambda_l}{2\sqrt{d-1}} \right). \end{aligned}$$

$$\text{Where : } \delta_T(x - y) = \frac{2}{\pi} \sum_{t=0}^{\infty} \frac{1}{1 + \delta_{t,0}} T_t(x) T_t(y)$$

$$\int_{-1}^1 \frac{dx}{\sqrt{1-x^2}} \delta_T(x - y) f(x) = f(y)$$

Note:

$$\text{tr } Y^0 = 2E \quad \rightarrow \quad y_0 = \frac{2E-1}{V} = d - \frac{1}{V}$$

$$\text{tr } Y^1 = \text{tr } Y^2 = 0 \quad \rightarrow \quad y_t = -\frac{1}{V} (d-1)^{\frac{t}{2}} \quad ; \quad t = 1, 2$$

## The trace formula:

$$\rho(\lambda) = \frac{d}{2\pi} \frac{\sqrt{4(d-1) - \lambda^2}}{d^2 - \lambda^2} + \frac{1}{\pi} \operatorname{Re} \left[ \sum_{t=3}^{\infty} \frac{y_t}{\sqrt{4(d-1) - \lambda^2}} e^{it \arccos \frac{\lambda}{2\sqrt{d-1}}} \right]$$

*Kesten McKay*

*Sum over t-periodic non-backscattering walks.*

$$y_t = \frac{1}{V} \frac{[\text{Number of } t \text{ periodic n.b. walks}] - [(d-1)^t - 1]}{(d-1)^{\frac{t}{2}}}.$$

$y_t$  : The deviation of the number of t-periodic non-backscattering walks from its mean. (properly normalized)

**Note:**  $\langle y_t \rangle_{\mathcal{G}} = 0$  for  $t \leq \log_{d-1} V$  (Bollobas, McKay, Wormald)

”Action” per single step on the orbit:  $\arccos \frac{\lambda}{2\sqrt{d-1}}$

Alternatively:

$$\rho(\lambda) = \frac{d}{2\pi} \frac{\sqrt{4(d-1) - \lambda^2}}{d^2 - \lambda^2} + \frac{1}{\pi} \sum_{t=3}^{\infty} \frac{y_t}{\sqrt{4(d-1) - \lambda^2}} T_t \left( \frac{\lambda}{2\sqrt{d-1}} \right)$$

$T_l(x)$  : Chebyshev Polynomials of the first kind.

$$\mathcal{R} = \{ \lambda : |\lambda| < 2(d-1)^{1/2} \}$$

**Non Ramanujan graphs :**  $\mathcal{R}^c$  not empty.

$$\text{tr } Y^t = (d-1)^t + (E-V)(1+(-1)^t) + 1 + 2(d-1)^{\frac{t}{2}} \left[ \sum_{\lambda_l \in \mathcal{R}} \cos t\phi_l + \sum_{\lambda_l \in \mathcal{R}^c} \cosh t\psi_m \right].$$

$$\psi_m = \text{arcosh} \frac{|\lambda_m|}{2\sqrt{d-1}}$$

$$\text{tr } Y^t = (d-1)^t + 2(d-1)^{\frac{t}{2}} \sum_{\lambda_l \in \mathcal{R}^c} \cosh t\psi_m + V \left( \frac{d}{2} - 1 \right) (1+(-1)^t) + 1 + 2(d-1)^{\frac{t}{2}} \sum_{\lambda_l \in \mathcal{R}} T_t \left( \frac{\lambda_l}{2\sqrt{d-1}} \right).$$

$$y_t = \frac{1}{V} \frac{[\text{Number of } t \text{ periodic n.b. walks}] - \left[ (d-1)^t \left( 1 + \frac{2}{(d-1)^{t/2}} \sum_{\lambda_l \in \mathcal{R}^c} \cosh t\psi_m \right) - 1 \right]}{(d-1)^{\frac{t}{2}}}.$$

$y_t$  : The deviation of the number of  $t$ -periodic non-backscattering walks from its exponentially growing part (properly normalized).

$$[\dots] \approx (d-1)^t \left( 1 + \exp \left[ -t \left( \frac{1}{2} \ln(d-1) - \left( \frac{V^{-0.6}}{\sqrt{d-1}} \right)^{\frac{1}{2}} \right) \right] \right) - 1$$

## 2.b A continuous family of trace formulas

Back to Bartholdi's identity :

$$\det(\eta I^{(2E)} - (B - wJ)) = (\eta^2 - w^2)^{E-V} \det(I^{(V)}(\eta^2 + w(d - w)) - \eta A).$$

The spectrum of  $Y(w) = B - wJ$  reads:

$$(d - w), w, (w, -w) \times (E - V), \left( \sqrt{w(d - w)} e^{\pm i\phi_k} \right)_{k=1, V-1}$$

For  $1 < w < d - 1$ ,  $2\sqrt{d - 1} < 2\sqrt{w(d - w)}$ .

Choose  $w$  such that  $\lambda_1 < 2\sqrt{w(d - w)} \rightarrow \phi_k = \arccos \frac{\lambda_k}{2\sqrt{w(d - w)}}$  are real.

**$Y(w)$  gives a weight  $(1-w)$  to back scattering!**

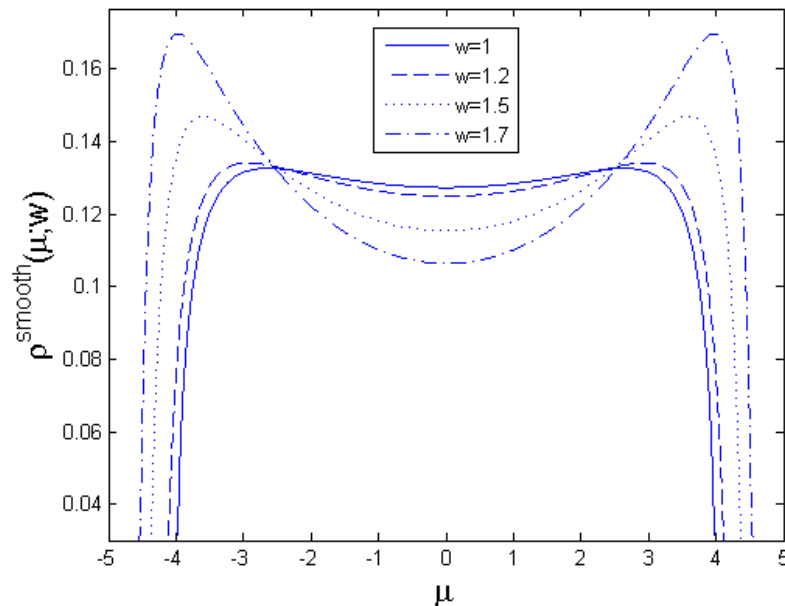
Define : 
$$y_t(w) = \frac{1}{V} \frac{\text{tr } Y(w)^t - ((d-w)^t + w^t)}{(w(d-w))^{\frac{t}{2}}}$$

**The  $w$  - trace formula:**

$$\rho(\lambda) = \rho^{\text{smooth}}(\lambda; w) + \frac{1}{\pi} \text{Re} \left[ \sum_{t=3}^{\infty} \frac{y_t(w)}{\sqrt{4w(d-w) - \lambda^2}} e^{it \frac{\lambda}{2\sqrt{w(d-w)}}} \right]$$

$$\rho^{\text{smooth}}(\lambda; w) = \frac{d}{2\pi \sqrt{4w(d-w) - \lambda^2}} \cdot \left( 1 - \frac{(d-2w)(d-2)}{d^2 - \lambda^2} + 2 \frac{(w-1)^2(2\lambda^2 - 1)}{w(d-w)} \right)$$

**The  $w$  - trace formula includes counting of backscattering walks weighted by  $(1-w)$  per backscatter!**



## 2.c The spectrum of equilateral quantum graphs from Bartholdi

For an equilateral  $d$ -regular quantum graph with  $L = \pi$

$$U = e^{ik\pi} \left( \frac{2}{d}B - J \right).$$

Taking  $\eta = \frac{d}{2}e^{-ik\pi}$ ,  $w = \frac{d}{2}$ , Bartholdi gives:

$$\det(I^{(2E)} - U(k)) = C(1 - e^{2ik\pi})^{E-V} \prod_{l=1}^V \left( \cos k\pi - \frac{\lambda_l}{d} \right)$$

$$\frac{1}{V}\rho_Q(k) = \frac{1}{2} \frac{\sqrt{4(d-1) - d^2 \cos^2(k\pi)}}{|\sin(k\pi)|} \cdot \mathbb{I}_R + \left( \frac{d-2}{2} \right) \delta(k-1)$$

$$\mathbb{I}_R = \begin{cases} 1 & \text{if } k \in \left[ \frac{1}{\pi} \arccos\left(\frac{2\sqrt{d-1}}{d}\right), 1 - \frac{1}{\pi} \arccos\left(\frac{2\sqrt{d-1}}{d}\right) \right] \\ 0 & \text{otherwise} \end{cases}.$$

The spectrum for the equilateral graph is periodic consisting of two components.

When the lengths are randomized  $l_{ij} = \pi(1 + \epsilon r_{ij})$  with  $|r_{ij}| \leq 1$  random uniformly distributed, the numerical distribution is:

