Other Results

QUANTUM DOTS AND OPTICAL CAVITIES

PHOTONS, COUPLED QUANTUM DOTS AND QUBITS

TWO EXCITONS IN QD WITH COHERENT FIELD





From the following article: Wiring up quantum systems R. J. Schoelkopf & S. M. Girvin Nature 451, 664-669(7 February 2008) doi:10.1038/451664a

EXITONIC MODEL

$$\begin{split} \widehat{H} &= \frac{1}{2} \hbar \omega \left[\hat{\sigma}_{+}, \hat{\sigma}_{-} \right] + \hbar \omega \, \hat{a}^{+} \hat{a} + i \hbar \lambda (\hat{\sigma}_{-} \hat{a}^{+} - \hat{\sigma}_{+} \hat{a}) \\ \\ \hat{\sigma}_{+} &= |e\rangle \langle g| \qquad \hat{\sigma}_{-} = |g\rangle \langle e| \end{split}$$

MASTER EQUATION

$$\begin{aligned} \frac{d\rho}{dt} &= -i[H,\rho] + L_{qd}^{A} + L_{qd}^{B} + L_{c} \\ L_{qd}^{A} &= \frac{g_{A}}{2} \left(2\sigma_{0}^{A}\rho\sigma_{1}^{A} - \sigma_{1}^{A}\sigma_{0}^{A}\rho - \rho\sigma_{1}^{A}\sigma_{0}^{A} \right) \\ L_{qd}^{B} &= \frac{g_{B}}{2} \left(2\sigma_{0}^{B}\rho\sigma_{1}^{B} - \sigma_{1}^{B}\sigma_{0}^{B}\rho - \rho\sigma_{1}^{B}\sigma_{0}^{B} \right) \\ L_{c} &= g_{c} \left(2a\rho a^{+} - a^{+}a\rho - \rho a^{+}a \right) \\ H &= \frac{\omega_{qd}^{A}}{2} [\sigma_{+}^{A}, \sigma_{-}^{A}] + \frac{\omega_{qd}^{B}}{2} [\sigma_{+}^{B}, \sigma_{-}^{B}] \\ &+ \omega_{c}a^{+}a + i\lambda^{A} \left(a^{+}\sigma_{-}^{A} - a\sigma_{+}^{A} \right) + i\lambda^{B} \left(a^{+}\sigma_{-}^{B} - a\sigma_{+}^{B} \right) \end{aligned}$$

RABI OSCLLIATIONS



> Time (T.I): 0.013ns

CORRELATION



SQUARE OF THE DENSITY TRACE OPERATOR



T.I: (0.001415ns -0. 12ns), peaks: 0.0046, 0.0066 y 0.0197ns

TWO EXCITONS AND SPIN OF QDs IN EMPTY FIELD

RABI OSCILLATIONS BY EXCITONS





DENSITY MATRIX

Table: Matrix Density of the QDs and photon in the cavity.

	$\langle 0 _{\scriptscriptstyle A} \langle 0 _{\scriptscriptstyle B} \langle 0 _{\scriptscriptstyle C}$	$\langle 0 _{_{\mathcal{A}}} \langle 0 _{_{\mathcal{B}}} \langle 1 _{_{\mathcal{C}}}$	$\langle 0 _{A} \langle 1 _{B} \langle 0 _{c}$	$\langle 1 _{\scriptscriptstyle A} \langle 0 _{\scriptscriptstyle B} \langle 0 _{\scriptscriptstyle c}$	
$\left 0\right\rangle_{A}\left 0\right\rangle_{B}\left 0\right\rangle_{c}$	a(t)	0	0	0	
$ 0\rangle_{_{\!\mathcal{A}}} 0\rangle_{_{\!\mathcal{B}}} 1\rangle_{_{\!\mathcal{C}}}$	0	b(t)	d(t)	d(t)	
$ 0\rangle_{A} 1\rangle_{B} 0\rangle_{c}$	0	d(t)	c(t)	c(t)	
$ 1\rangle_{_{\!A}} 0\rangle_{_{\!B}} 0\rangle_{_{\!c}}$	0	d(t)	c(t)	c(t)	

	$\langle \downarrow _A \langle \downarrow _B \langle 0 _c \langle 0 _L$	${\langle \downarrow }_{A} {\langle \downarrow }_{B} {\langle 1 }_{c} {\langle 0 }_{L}$	$\left<\uparrow\right _A\left<\downarrow\right _B\left<0\right _c\left<1\right _L$	${\langle \downarrow }_{A}{\langle \uparrow }_{B}{\langle 0 }_{c}{\langle 1 }_{L}$
$ \downarrow\rangle_A \downarrow\rangle_B 0\rangle_c 0\rangle_L$	a(t)	0	0	0
$ {\downarrow}\rangle_{A} {\downarrow}\rangle_{B} 1\rangle_{c} 0\rangle_{L}$	0	b(t)	d(t)	d(t)
$ \!\!\uparrow\rangle_{\!_{A}} \!\!\downarrow\rangle_{\!_{B}} 0\rangle_{c} 1\rangle_{\!_{L}}$	0	d(t)	c(t)	c(t)
$ {\downarrow}\rangle_{_{A}} {\uparrow}\rangle_{_{B}} 0\rangle_{_{c}} 1\rangle_{_{L}}$	0	d(t)	c(t)	c(t)

EVOLUTION OF d(t)



DENSITY MATRIX DIAGONALIZATION

Eigenvectors:

$$|\phi_1\rangle = |0\rangle_A |0\rangle_B |0\rangle_c$$

$$|\phi_2\rangle = \frac{\pm \sqrt{b(t)}|0\rangle_A |0\rangle_B |1\rangle_c + \sqrt{c(t)} (|0\rangle_A |1\rangle_B |0\rangle_c + |1\rangle_A |0\rangle_B |0\rangle_c)$$

$$\sqrt{b(t) + 2c(t)}$$

	$\langle \phi_1 $	$\langle \phi_2 $
$ \phi_{\rm i} angle$	a(t)	0
$ \phi_2 angle$	0	b(t) + 2c(t)

EVOLUTION OF THE ENTANGLEMENT STATES FOR EXCITONS



T.I: 0.0286ns

EVOLUTION OF THE ENTANGLEMENT STATES WITH SPIN



≻T.I: 0.0715ns

TOTAL ENTROPY

 $S(\rho_{A,B,c}) = -(a(t)\ln[a(t)] + (b(t) + 2c(t))\ln[b(t) + 2c(t)]) / \ln[2]$



DIAGONALIZATION OF RESTRICTED DENSITY OPERATOR IN A AND B

Eigenvectors:

$$|\Phi_{1}\rangle = |0\rangle_{A}|0\rangle_{B}$$
$$|\Phi_{2}\rangle = \frac{1}{\sqrt{2}} (|0\rangle_{A}|1\rangle_{B} + |1\rangle_{A}|0\rangle_{B})$$

	$\langle \Phi_1 $	$\langle \Phi_2 $
$ \Phi_1 angle$	a(t)+b(t)	0
$ \Phi_2 angle$	0	2c(t)

EVOLUTION OF THE ENTANGLEMENT STATES OF REDUCED EXCITONS





SQUARE TRACE OF DENSITY OPERATOR BY EXCITONS



Grado de mezcla y entrelazamiento



CORRELATION OF EXCITONS





ENVIRONMENTS NO DISSIPATIVE

Equations to solve:

 $\frac{d(a(t))}{dt} = 0 \frac{d(b(t))}{dt} = -4\lambda d(t) \quad \frac{d(c(t))}{dt} = 2\lambda d(t) \quad \frac{d(d(t))}{dt} = \lambda(b(t) - 2c(t))$ $\succ \text{Solutions:}$

a(t) = 0 $b(t) = \cos^{2}(\sqrt{2\lambda}t)$ $c(t) = \frac{1}{2}\sin^{2}(\sqrt{2\lambda}t)$ $d(t) = \frac{1}{\sqrt{2}}\sin(\sqrt{2\lambda}t)\cos(\sqrt{2\lambda}t)$

MATRIX DENSITY

	(o[_(o]_(a]_	(0]_(1]_(0]_	(i], (o], (o],
$ 0\rangle_{s} 0\rangle_{s} 1\rangle_{s}$	b(t)	d(t)	d(t)
0 <u>),</u> 1 <u>),</u> 0 <u>)</u> ,	d(t)	c(t)	c(t)
$ 1\rangle_{s} 0\rangle_{s} 0\rangle_{s}$	d(t)	c(t)	c(t)

► EIGENVECTOR:

$$|\phi_{2}\rangle = \frac{\pm \sqrt{b(t)}|0\rangle_{A}|0\rangle_{B}|1\rangle_{c} + \sqrt{c(t)}(|0\rangle_{A}|1\rangle_{B}|0\rangle_{c} + |1\rangle_{A}|0\rangle_{B}|0\rangle_{c}}{\sqrt{b(t)+2c(t)}}$$

DECOHERENCE IN EXCITONS





DIAGONALIZATION OF REDUCED DENSITY OPERATOR IN A AND B

Eigenvectors:

$$|\Phi_1\rangle = |0\rangle_A|0\rangle_B$$

$$\left|\Phi_{2}\right\rangle = \frac{1}{\sqrt{2}} \left(\left|0\right\rangle_{A}\left|1\right\rangle_{B} + \left|1\right\rangle_{A}\left|0\right\rangle_{B}\right)$$

	$\langle \Phi_1 $	$\langle \Phi_2 $
$ \Phi_1 angle$	b(t)	0
$ \Phi_2 angle$	0	2c(t)

TRACE OF SQUARE OF DENSITY OPERATOR



RABI FREQUENCES

\geq EXCITONES= $\sqrt{2} \times \lambda = \sqrt{2} \times 315$ GHz

\geq ESPIN= $\sqrt{2^*\lambda_{eff}}$ = $\sqrt{2^*}$ 24.18GHz

ONE EXCTION OR SPIN

$$\frac{d(b(t))}{dt} = -2\lambda d(t)$$
$$\frac{d(c(t))}{dt} = 2\lambda d(t)$$
$$\frac{d(d(t))}{dt} = \lambda(b(t) - c(t))$$

	$ \downarrow\rangle_{QD} 1\rangle_{c}$	$ \uparrow \rangle_{QD} 0 \rangle_c$
$ \downarrow\rangle_{QD} 1\rangle_{c}$	b(t)	d(t)
$ \uparrow\rangle_{QD} 0\rangle_{c}$	d(t)	c(t)

$$b(t) = \cos^{2}(\lambda t)$$
$$c(t) = \sin^{2}(\lambda t)$$
$$d(t) = \sin(\lambda t) * \cos(\lambda t)$$

CONCLUSIONS

Excitons interaction in the quantum dot with a coherent field, the interlaced state is not defined, but the range of entanglement was predominant during the system dynamics.

In the interaction of quantum dot excitons with empty field in times proportional to a half-integer number of π on Rabi frequency were obtained maximally entangled states as Bell states, useful in computer science and information quantum. In the interaction of quantum dot spins with empty field, the dynamics were similar to that of excitons with empty field, but in this model, the frequency of Rabbi and coherence times are greater because model conditions. ➢ Spin model, is predominant over the exciton due to the coherence time exceeds the time for a certain computation operation (0.04ns). They analyze the dynamics of a single quantum dot (with exciton or spin) interacting with the different fields in the cavity, we obtained results that analyze the behavior of quantum gates with two-level systems

RECTANGULAR DOUBLE BARRIER POTENTIAL



$$A_{\sigma,\sigma'} = \frac{e\hbar}{2m^*} \vec{\sigma} \cdot \vec{B}$$

$$\langle I \rangle = \frac{2eT}{\hbar} \int_{-\infty}^{+\infty} d\omega [G_{01}^{+-}(\omega) - G_{10}^{+-}(\omega)]$$

$$I = 16\pi^3 J_0(B) T^2 \sum_{n=1}^{N} \int_{(n+\frac{1}{2})\hbar\omega}^{E_F} \frac{\rho_L(\hbar\omega)\rho_R(\hbar\omega)d(\hbar\omega)}{|\Lambda(\hbar\omega)|^2}$$

 $|\Lambda(\hbar\omega)|^2 = (1 - g^a_{LL} g^a_{RR} V^2) (1 - g^{\gamma}_{LL} g^{\gamma}_{RR} V^2)$











JJJ₀(B)













PÖSCHL-TELLER DOUBLE BARRIER POTENTIAL

RASHBA AND DRESSELHAUS EFFECTS



Anticrossing due to Rashba coupling



 $\mu \Phi$