

Decoherence and relaxation of quantum systems

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Thesis

- Part I: Decoherence and relaxation in an open system
 - Bath of non-interacting electrons *Restrepo, Camalet, Chitra, Dupont Phys Rev. B(2011)*
 - Bath with a long range order *Restrepo, Camalet, Chitra
(to be submitted)*
- Part II: Thermalization in a closed system *Restrepo, Camalet
New J. Phys (2010)*

Thesis

- Part I: Decoherence and relaxation in an open system
 - Bath of non-interacting electrons *Restrepo, Camalet, Chitra, Dupont Accepted Phys Rev. B*
 - Bath with a long range order *Restrepo, Camalet, Chitra (in preparation)*
- Part II: Thermalization in a closed system *Restrepo, Camalet New J. Phys (2010)*

Decoherence and relaxation of a qubit coupled to a bath with long range order

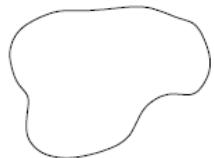
Qubit + bath



- Qubit $|\psi\rangle = \alpha |\uparrow\rangle + \beta |\downarrow\rangle$

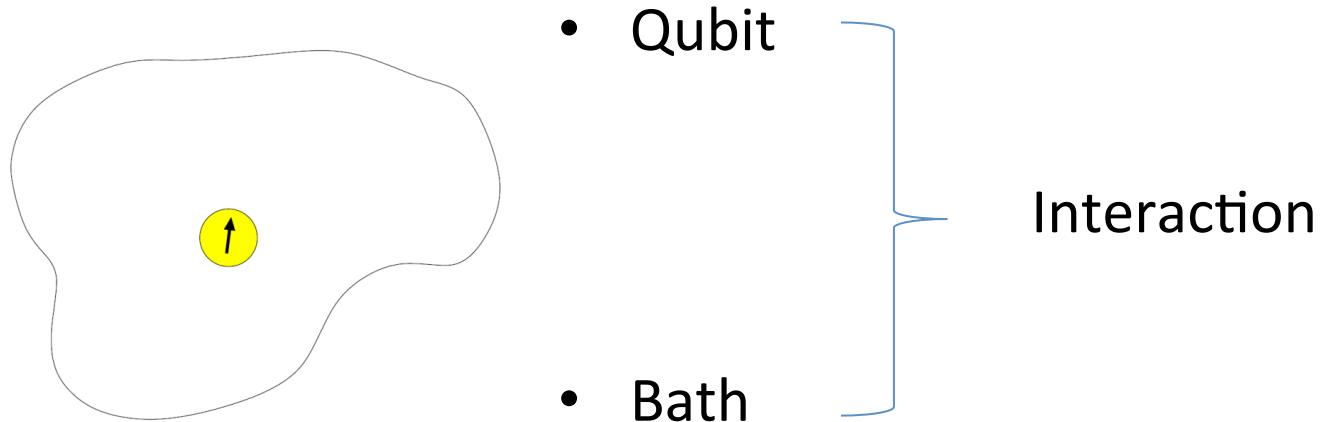
$$\rho_s = |\psi\rangle \langle\psi| = \begin{bmatrix} |\alpha|^2 & \alpha\beta^* \\ \alpha^*\beta & |\beta|^2 \end{bmatrix}$$

Density matrix



- Bath ρ_B

Decoherence and relaxation



$$\rho_s = \begin{bmatrix} |\alpha|^2 & \alpha\beta^* \\ \alpha^*\beta & |\beta|^2 \end{bmatrix}$$

Time evolution

- Decoherence
- Relaxation

Why study decoherence and relaxation?

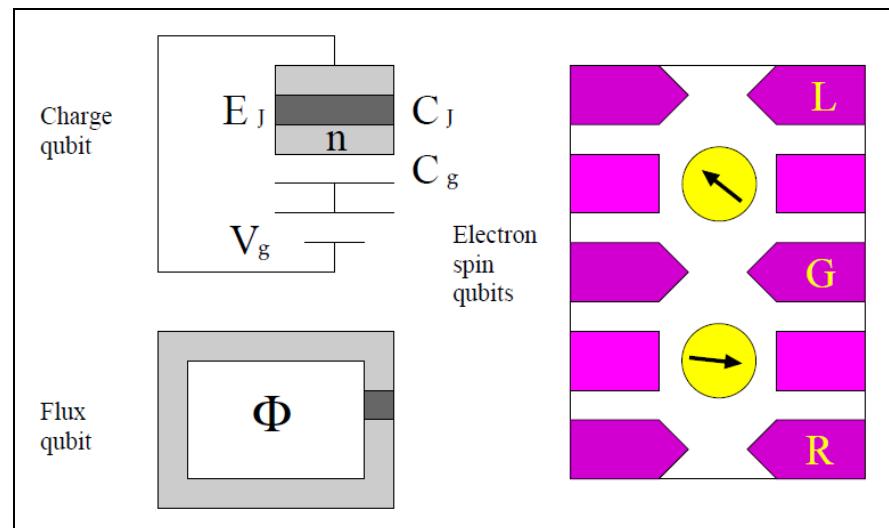
- For quantum computation it is important to have coherent qubits.
- Qubits can be used to probe the environment.

Physical realizations of qubits

Solid state qubits → Size advantage

Examples

- Josephson junctions
- Spin in quantum dots



*Makhlin
Rev Mod Phys (2001)*

*DiVincenzo, Loss
Phys Rev A (1998)*

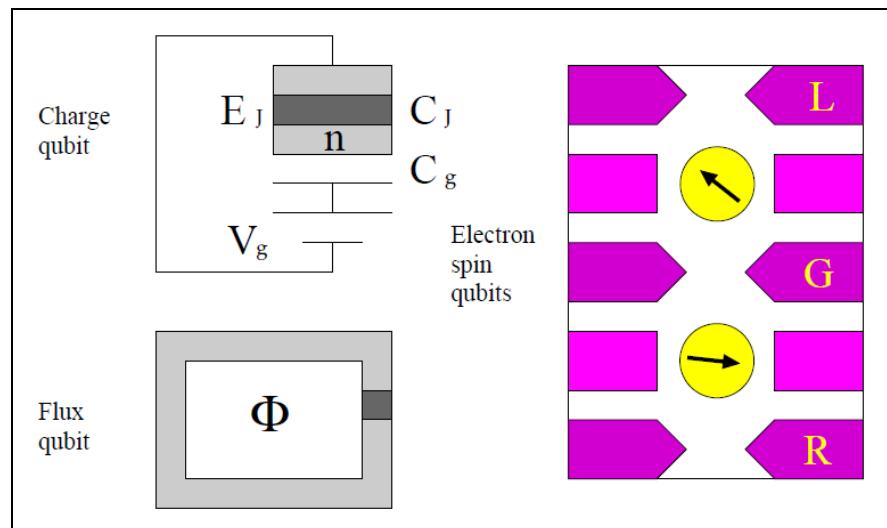
Physical realizations of qubits

Solid state qubits → Size advantage

Examples

- Josephson junctions
- Spin in quantum dots

→ Interacting nuclear spins
and electronic baths



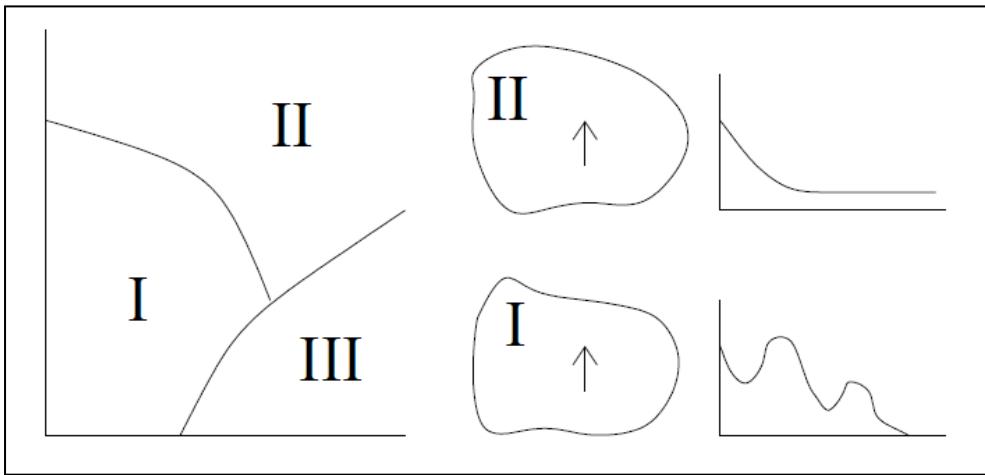
*Coish, Loss
Phys Rev B (2008)*

*Glazman, Loss
Phys Rev Lett (2002)*

*Schiller et al
Phys Rev B (2006)*

*Yamada et al
Conf Proc (2007)*

Qubits as probes



Phase diagram

Decoherence

- Theoretical

*Chitra, Camalet
Phys Rev Lett (2007)*

*Winograd , Chitra,
Rozenberg Phys Rev B
(2010)*

*Vernier / Jiang
Phys Rev A (2011)*

- Experimental

*Buttiker et al
Phys Rev B (2010)*

*Vandersypen et al
Rev Mod Phys (2005)*

State of art

Baths	• Nuclear spins	<i>Glazman, Loss</i> <i>Phys Rev Lett</i> (2002)	<i>Coish, Loss</i> <i>Phys Rev B</i> (2008)
	• Electrons	<i>Schiller et al</i> <i>Phys Rev B</i> (2006)	<i>Yamada et al</i> <i>Conf Proc</i> (2007)
Effect of	• Bath Interactions	<i>Paganelli et al</i> <i>Phys Rev A</i> (2002)	<i>Tessieri et al</i> <i>Jour Phys A</i> (2003)
	• Phase transition	<i>Yuan et al</i> <i>EPL</i> (2005)	<i>Chitra, Camalet</i> <i>Phys Rev Lett</i> (2007)
		<i>Wang et al</i> <i>Phys Lett A</i> (2008)	

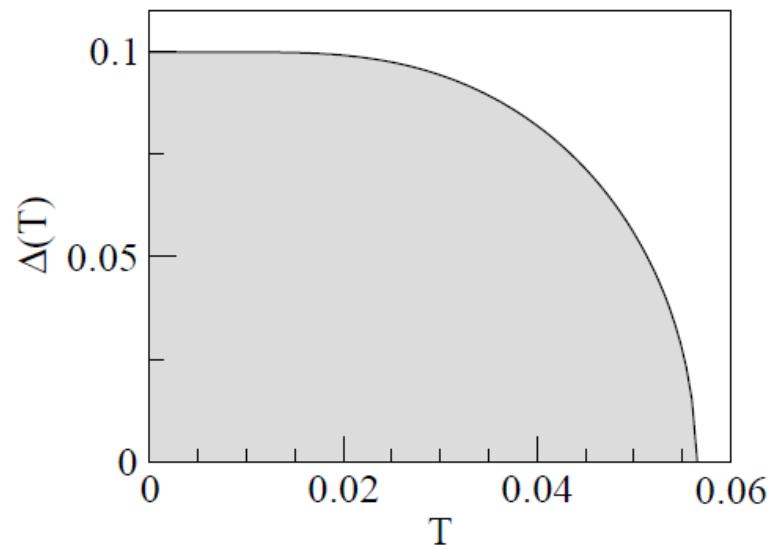
What is the effect of order in the bath on decoherence and relaxation of the qubit?

Model

$$H = H_S + H_B + H_I$$

- The **qubit** has no intrinsic dynamics
- The **bath** is described by BCS Hamiltonian
 - $T < T_c$ superconductor
 - $T > T_c$ metal
- Interaction

$$H_S = 0$$



$$H_I = \sum_{\beta=0,x,y,z} V_\beta \sigma_\beta^c$$

Interaction

- Charge

$$H_I = \lambda \sigma_z^c \sum_{kp} \left(c_{k\uparrow}^\dagger c_{p\uparrow} + c_{k\downarrow}^\dagger c_{p\downarrow} \right)$$

- Kondo

$$H_I = \lambda \boldsymbol{\sigma}^c \cdot \sum_{k,p,\alpha\beta} c_{k\alpha}^\dagger \boldsymbol{\sigma}_{\alpha\beta} c_{p\beta}$$

- Order

$$H_I = \lambda \sigma_-^c \sum_{kp} c_{k\uparrow}^\dagger c_{p\downarrow}^\dagger - \lambda \sigma_+^c \sum_{kp} c_{k\uparrow} c_{p\downarrow}$$

Reduced dynamics of the qubit

- Initial state

$$\rho(0) = \rho_s(0) \otimes \rho_B(0)$$

$$\rho_s(0) = \begin{bmatrix} |\alpha|^2 & \alpha\beta^* \\ \alpha^*\beta & |\beta|^2 \end{bmatrix}$$

$$\rho_B(0) = \frac{e^{-H_\beta/T}}{Z_B}$$

Reduced dynamics of the qubit

- Initial state
- Unitary evolution
- Reduced density matrix

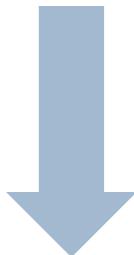
$$\rho(0) = \rho_s(0) \otimes \rho_B(0)$$

$$\frac{d}{dt} \rho = -i [H, \rho]$$

$$\rho_s(t) = \text{Tr}_B [\rho(t)]$$

Weak coupling techniques

Master equation $\frac{d}{dt}\rho_s(t) = \int dt' \Sigma(t', t)\rho_s(t')$



Born approximation λ^2

$$\frac{d}{dt}\rho_s(t) = \int dt' \Sigma^{(2)}(t', t)\rho_s(t')$$

Nakajima-Zwanzig (NZ)



Local approximation

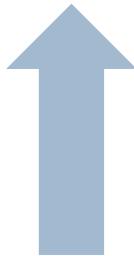
$$\frac{d}{dt}\rho_s(t) = \int dt' \Sigma^{(2)}(t', t)\rho_s(t')$$

Time convolutionless (TCL)

Master equation in Born approximation

Laplace transform $\rho_s(z) = \frac{1}{2} \sum_{\beta=0,x,y,z} M_\beta(z) \sigma_\beta^c$

$$z M_\beta(z) - \sum_{\alpha} h_{\beta\alpha} M_\alpha(z) - \sum_{\alpha} \Sigma_{\beta\alpha}(z) M_\alpha(z) = \langle \sigma_\beta^c \rangle_0$$



First order

$$\lambda \langle V_\gamma \rangle$$



Second order

$$\lambda^2 \langle V_\gamma(t) V_\delta \rangle$$

Self energies → Time evolution

Kondo coupling

- The problem is spin isotropic
- Bath operators have no net moment $\textcolor{brown}{h} = 0$



Only one self energy $\Sigma(z)$

Kondo coupling

- The problem is spin isotropic
- Bath operators have no net moment



Relaxation and decoherence are the same

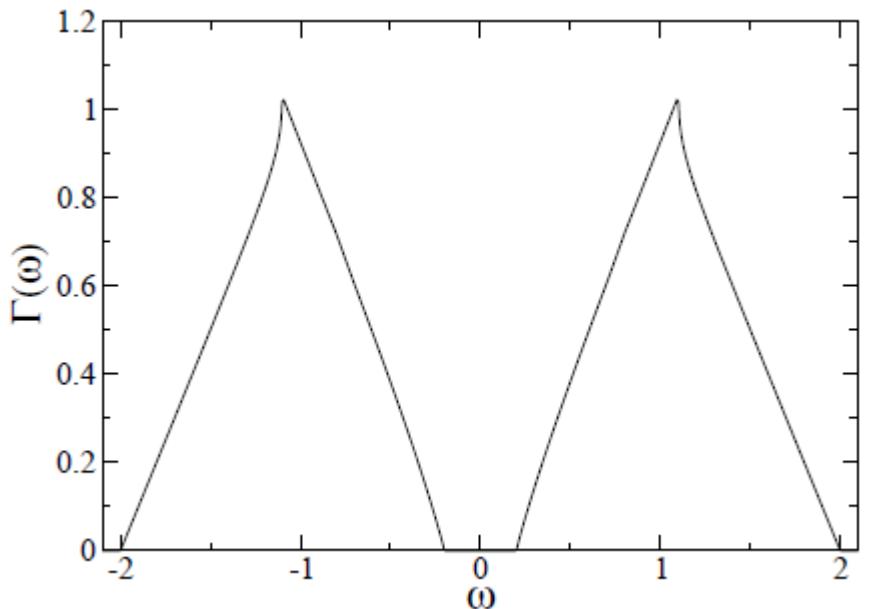
$$M_\beta(z) = [z - \underbrace{\Sigma(z)}_{M(z)}]^{-1} \langle \sigma_\beta^c \rangle_0$$

Dynamics in NZ and TCL

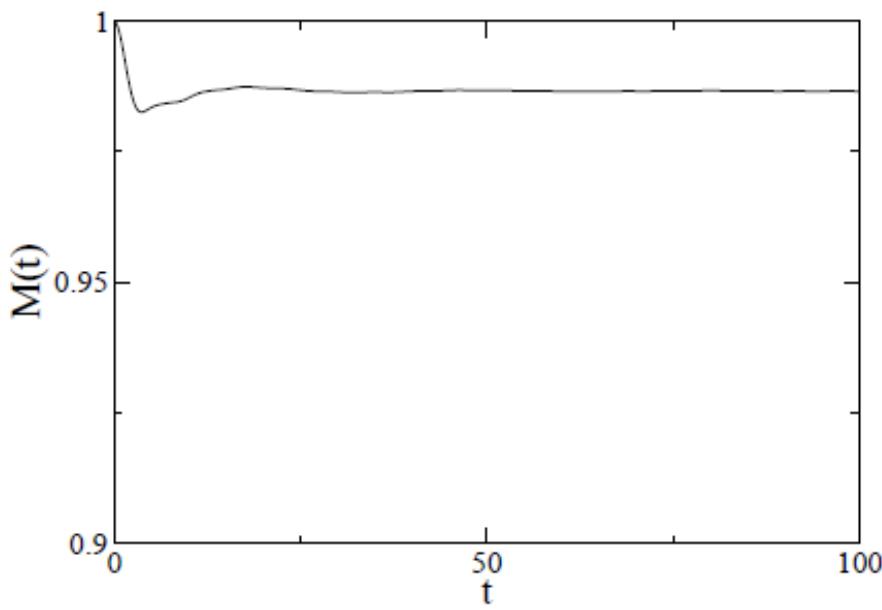
$$\Gamma(\omega) = -\Im m \lim_{\eta \rightarrow 0} \Sigma(\omega + i\eta)$$

- Time convolutionless $\ln M_{TCL}(t) = -\frac{2}{\pi} \int d\omega \frac{\sin^2 \omega t/2}{\omega^2} \Gamma(\omega)$
 - Nakajima-Zwanzig $M_{NZ}(t) = \frac{1}{\pi} \int_0^\infty d\omega \cos(\omega t) \tilde{\Gamma}(\omega)$
- Markovian asymptotic evolution $M(t) \simeq e^{-\Gamma(0)t}$  $\Gamma(\omega)$

Asymptotic dynamics for T=0 (TCL, NZ)



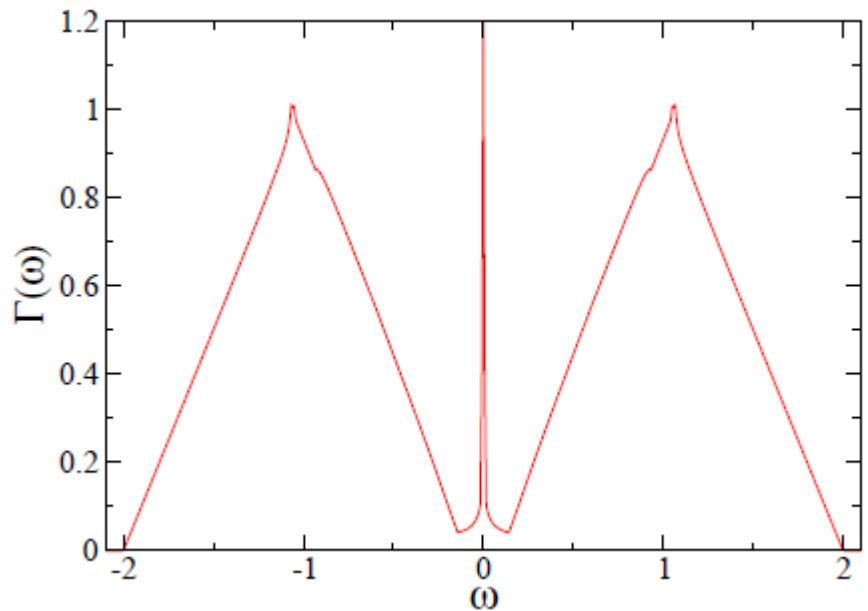
Gap at $\omega=0$



Incomplete decoherence/
relaxation

$M(t) \xrightarrow{t \rightarrow \infty} \text{constant}$

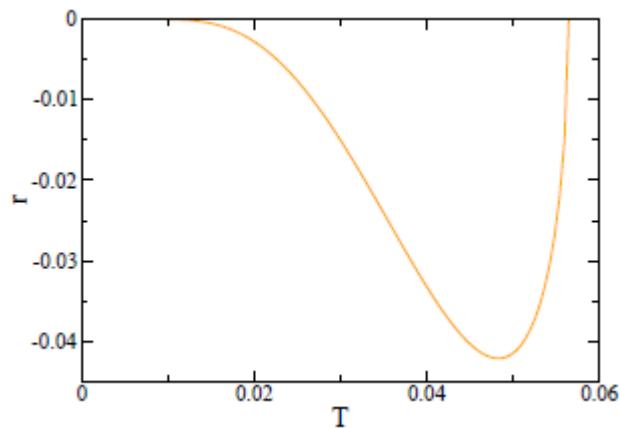
Self energy at finite temperature



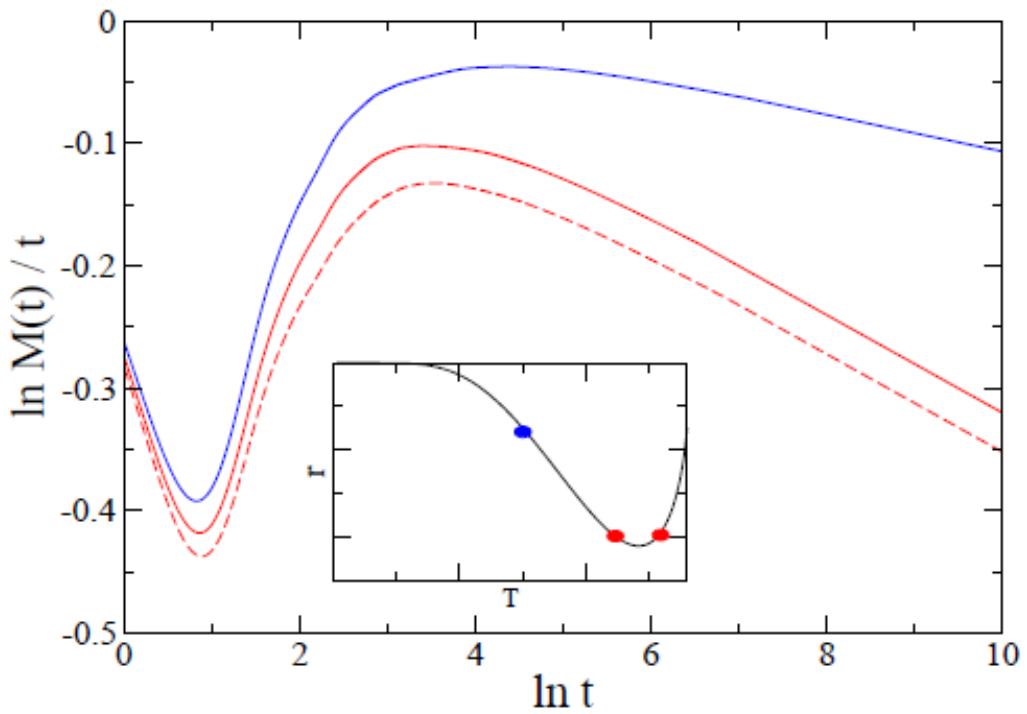
r is non monotonic

Logarithmic divergence at $\omega=0$

$$\frac{\Gamma(\omega)}{2\pi\lambda^2} = r \ln \frac{\omega}{T}$$



Asymptotic dynamics for $0 < T < T_c$ (TCL)



$$M_{TCL}(t) \simeq t^{2\pi\lambda^2 r(T)} t$$

Anomalous decoherence

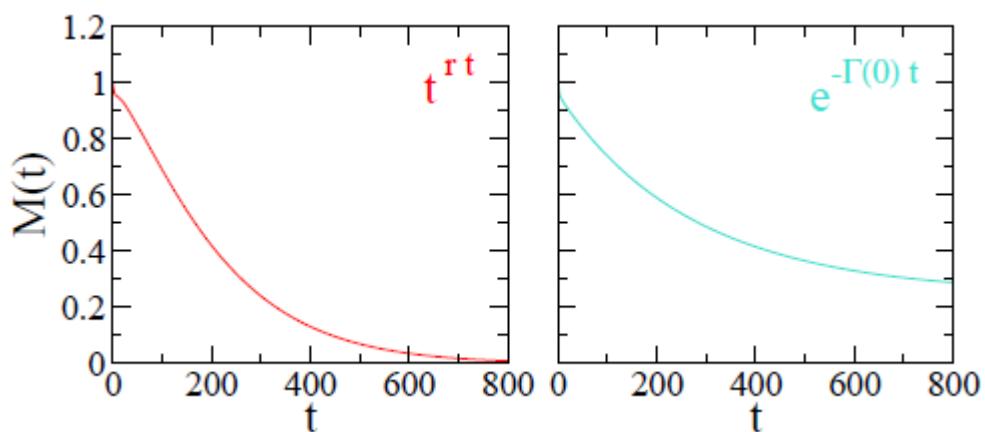
- Non Markovian
- Ultrafast
- “Reentrance”

Scales

- Small T ($t \gg 1/T$)
- Large T ($t \gg 1/\Delta$)

Asymptotic dynamics at $0 < T < T_c$ (TCL)

$$M_{TCL}(t) \simeq t^{2\pi\lambda^2 r(T)t}$$



Anomalous decoherence

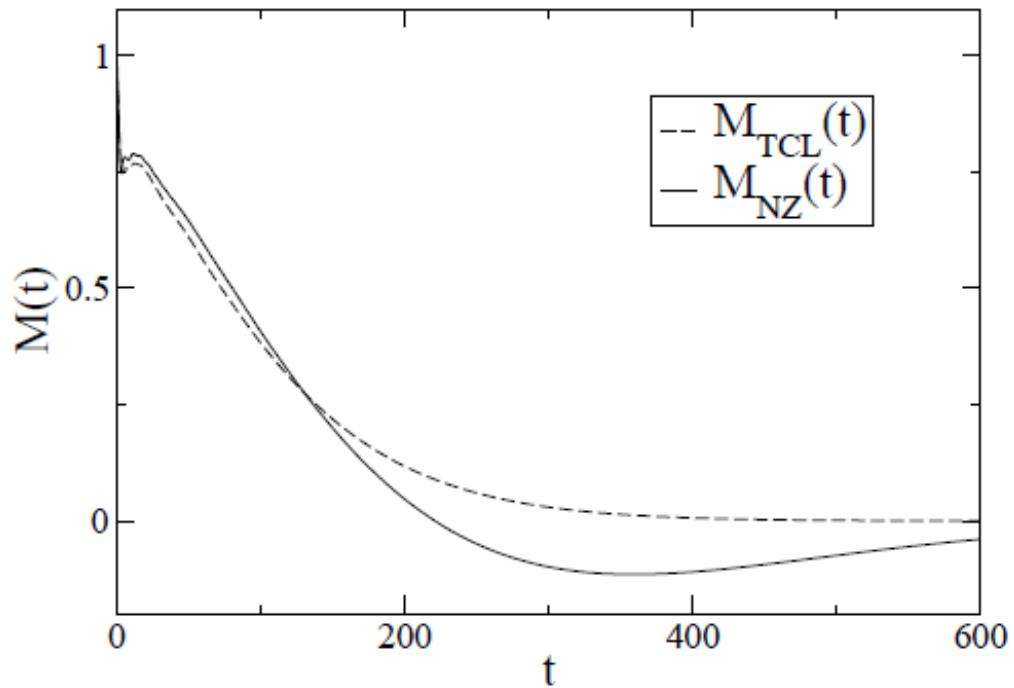
- Non Markovian
- Ultrafast
- Reentrance

Scales

- Small T ($t \gg 1/T$)
- Large T ($t \gg 1/\Delta$)

Asymptotic dynamics for $0 < T < T_c$ (NZ)

$$M_{NZ}(t) \simeq -\frac{1}{2\pi^2 \lambda^2 r} \frac{1}{t \ln t}$$



- Non Markovian decoherence/relaxation
- Different from TCL

Summary for Kondo coupling

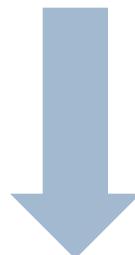
- Relaxation = Decoherence
- At $T=0$ incomplete decoherence (TCL and NZ)
- For $0 < T < T_c$ non Markovian decoherence ($TCL \neq NZ$)
- For $0 < T < T_c$ (TCL) Ultrafast non Markovian
→ disastrous for qubits

signature of order?

Order coupling

$$H_I = \lambda \sigma_-^c \sum_{kp} c_{k\uparrow}^\dagger c_{p\downarrow}^\dagger - \lambda \sigma_+^c \sum_{kp} c_{k\uparrow} c_{p\downarrow}$$

- The problem is not isotropic
- Bath operators have net moment $h_{yz} \propto \Delta$

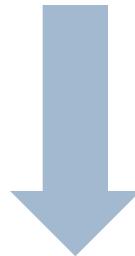


Self energy matrix

Relaxation \neq Decoherence

Order coupling

- The problem is not isotropic
- Bath operators have net moment $h_{yz} \propto \Delta$

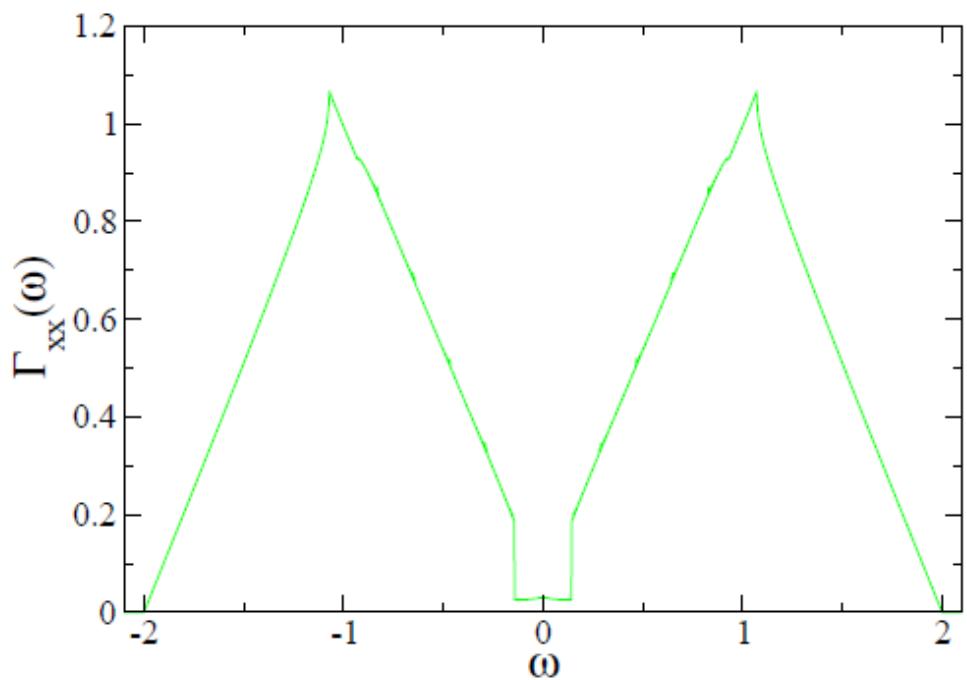


Self energy matrix

$$\begin{pmatrix} M_0(z) \\ M_x(z) \\ M_y(z) \\ M_z(z) \end{pmatrix} \begin{pmatrix} z & 0 & 0 & 0 \\ 0 & z - \Sigma_{xx} & 0 & 0 \\ 0 & 0 & z - \Sigma_{yy} & -h_{yz} \\ 0 & 0 & -h_{zy} & z - \Sigma_{zz} \end{pmatrix} = \begin{pmatrix} \langle \sigma_0^c \rangle_0 \\ \langle \sigma_x^c \rangle_0 \\ \langle \sigma_y^c \rangle_0 \\ \langle \sigma_z^c \rangle_0 \end{pmatrix}$$

The x component for $0 < T < T_c$ (TCL)

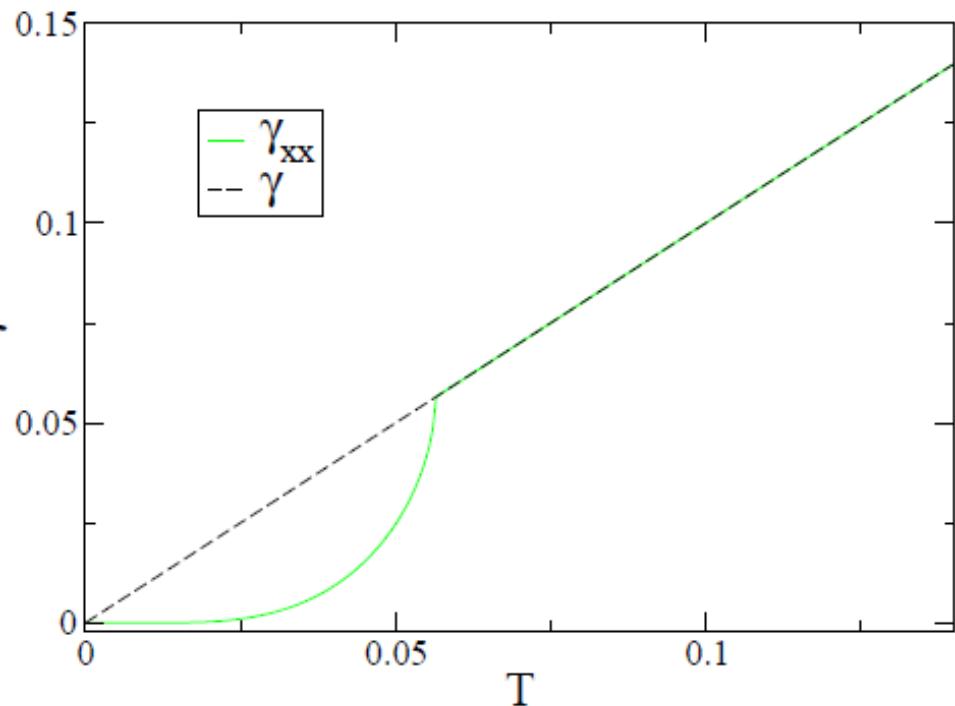
$$\Gamma_{xx}(\omega) \longrightarrow M_x(z)$$
$$\langle \sigma_x^c(t) \rangle$$



- Related to charge fluctuations
- Finite at $\omega=0 \rightarrow$ Markovian evolution

The x component for $0 < T < T_c$ (TCL)

$$\Gamma_{xx}(\omega) \longrightarrow M_x(z)$$
$$\langle \sigma_x^c(t) \rangle$$

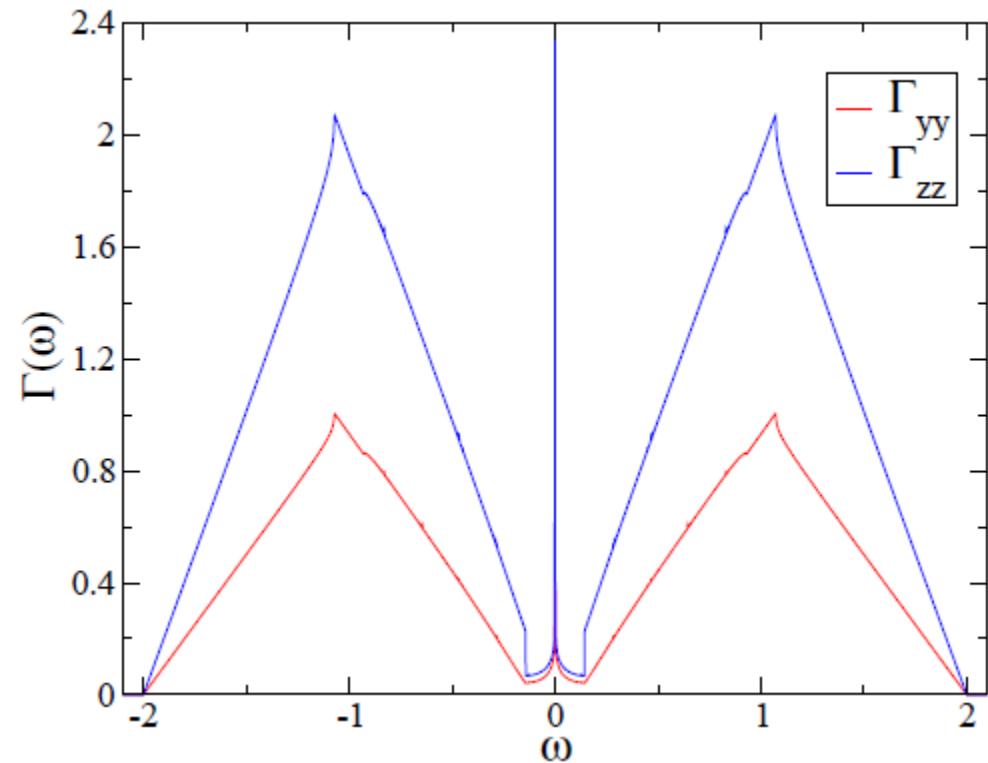


- Markovian
- Very slow

*Black and Fulde
Jour Phys Lett (1979)*

The y and z components for $0 < T < T_c$ (TCL)

$$h_{yz} \quad \Gamma_{zz}(\omega) \quad \Gamma_{yy}(\omega) \quad \longrightarrow \quad M_y(z) \quad M_z(z) \\ \langle \sigma_y^c(t) \rangle \quad \langle \sigma_z^c(t) \rangle$$



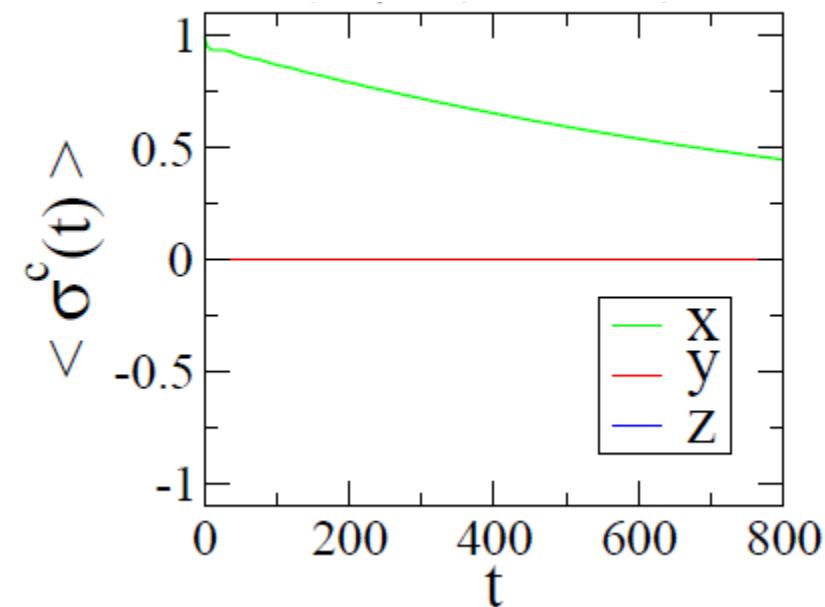
- Related to spin fluctuations
- Infrared divergence → non Markovian ultrafast evolution
- First order term → oscillations

Sensitivity to initial conditions

$$\begin{pmatrix} M_0(z) \\ M_x(z) \\ M_y(z) \\ M_z(z) \end{pmatrix} \begin{pmatrix} z & 0 & 0 & 0 \\ 0 & z - \Sigma_{xx} & 0 & 0 \\ 0 & 0 & z - \Sigma_{yy} & -h_{yz} \\ 0 & 0 & -h_{zy} & z - \Sigma_{zz} \end{pmatrix} = \begin{pmatrix} \langle \sigma_0^c \rangle_0 \\ \langle \sigma_x^c \rangle_0 \\ \langle \sigma_y^c \rangle_0 \\ \langle \sigma_z^c \rangle_0 \end{pmatrix}$$

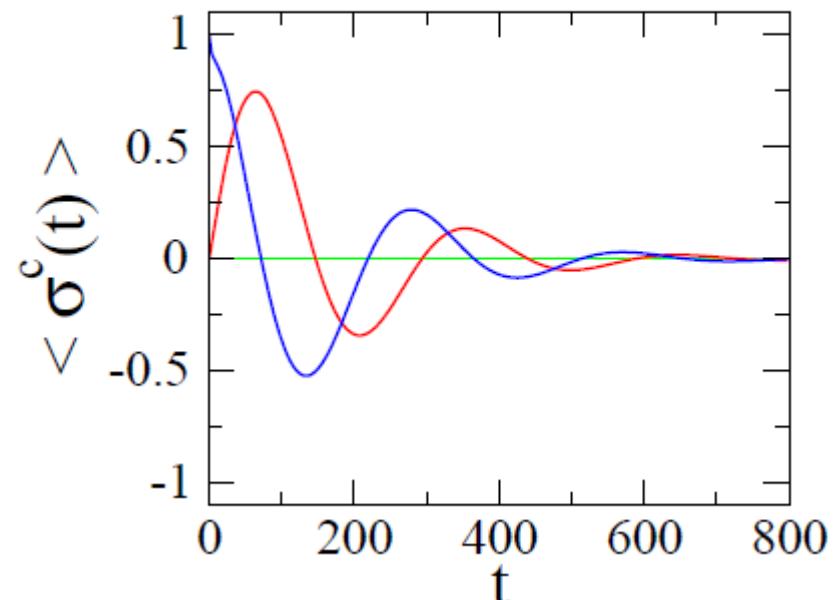
Sensitivity to initial conditions

$$|\psi\rangle = \frac{1}{\sqrt{2}} |\uparrow\rangle + \frac{1}{\sqrt{2}} |\downarrow\rangle$$



- No relaxation
- Markovian decoherence

$$|\psi\rangle = |\uparrow\rangle$$



- Non Markovian relaxation
- Non Markovian decoherence

Summary for Order coupling

- Relaxation \neq Decoherence
 - Sensitivity to initial conditions
 - Difference between x, y, z
 - Charge channel: Markovian and slow
 - Spin channel: Like Kondo with oscillations
- Important to couple to « good operators » and measure « good quantities »
- disastrous or good for qubits

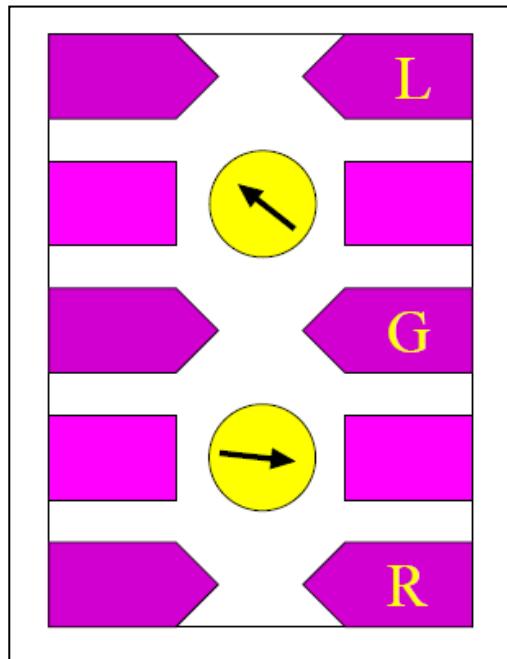
Perspectives

- $T \rightarrow T_c$ from metallic side
- Next order in TCL, NZ
- Other ordered baths
- Two qubits (entanglement and linear entropy)
- One qubit coupled to two baths:
 - Order + bosonic
 - Two ordered at different temperatures
- Strong coupling (Numerical approach)

Thank you, Gracias

P1: Dos qubits

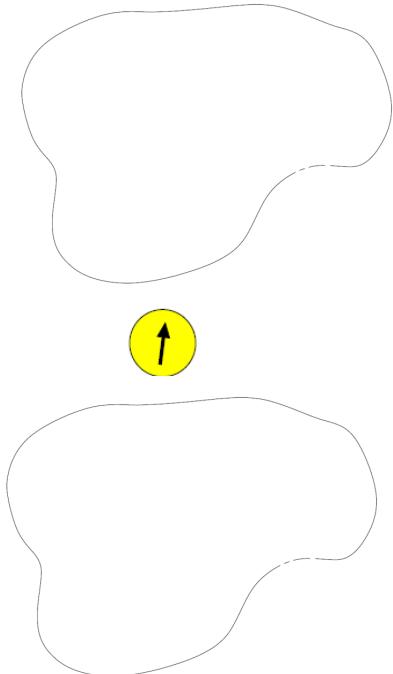
¿Cuál es la dinámica de dos qubits cuando interactúan con un baño en fase ordenada?



- Importante para aplicaciones, experimentos
- Entrelazamiento, entropía lineal
- Extensión « facil »

P2: Dos baños en equilibrio

¿Cuál es la decoherencia de un qubit cuando interactua con un baño de bosones y uno ordenado?

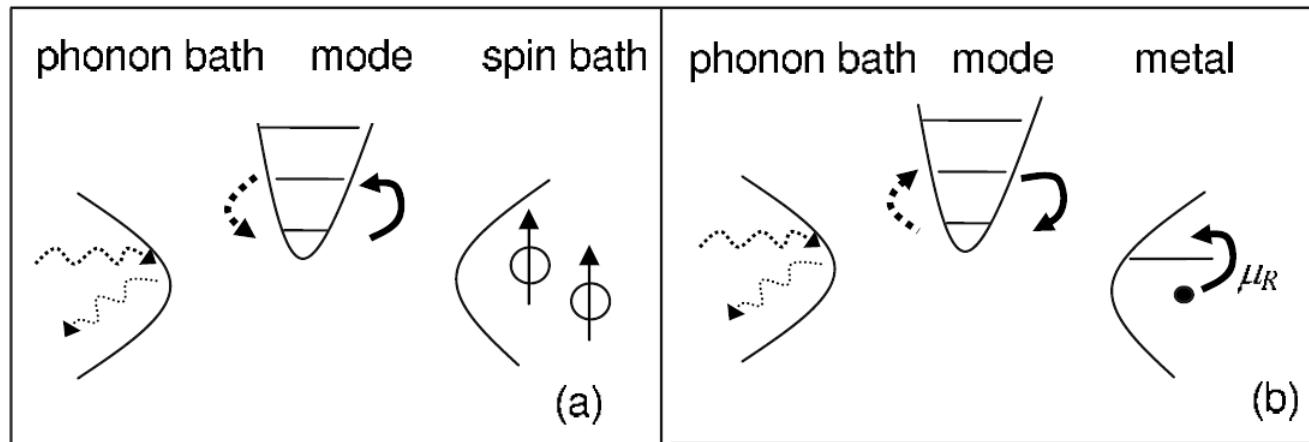


- Modelo más real
- Cada baño -> dinámicas diferentes
- ¿Quién gana?

R. Chitra (ETH, Zurich) S. Camalet (LPTMC, Paris)

P3: Dos baños a diferente temperatura

¿Cuál es la decoherencia de un qubit cuando interactua con dos baños a diferente temperatura?



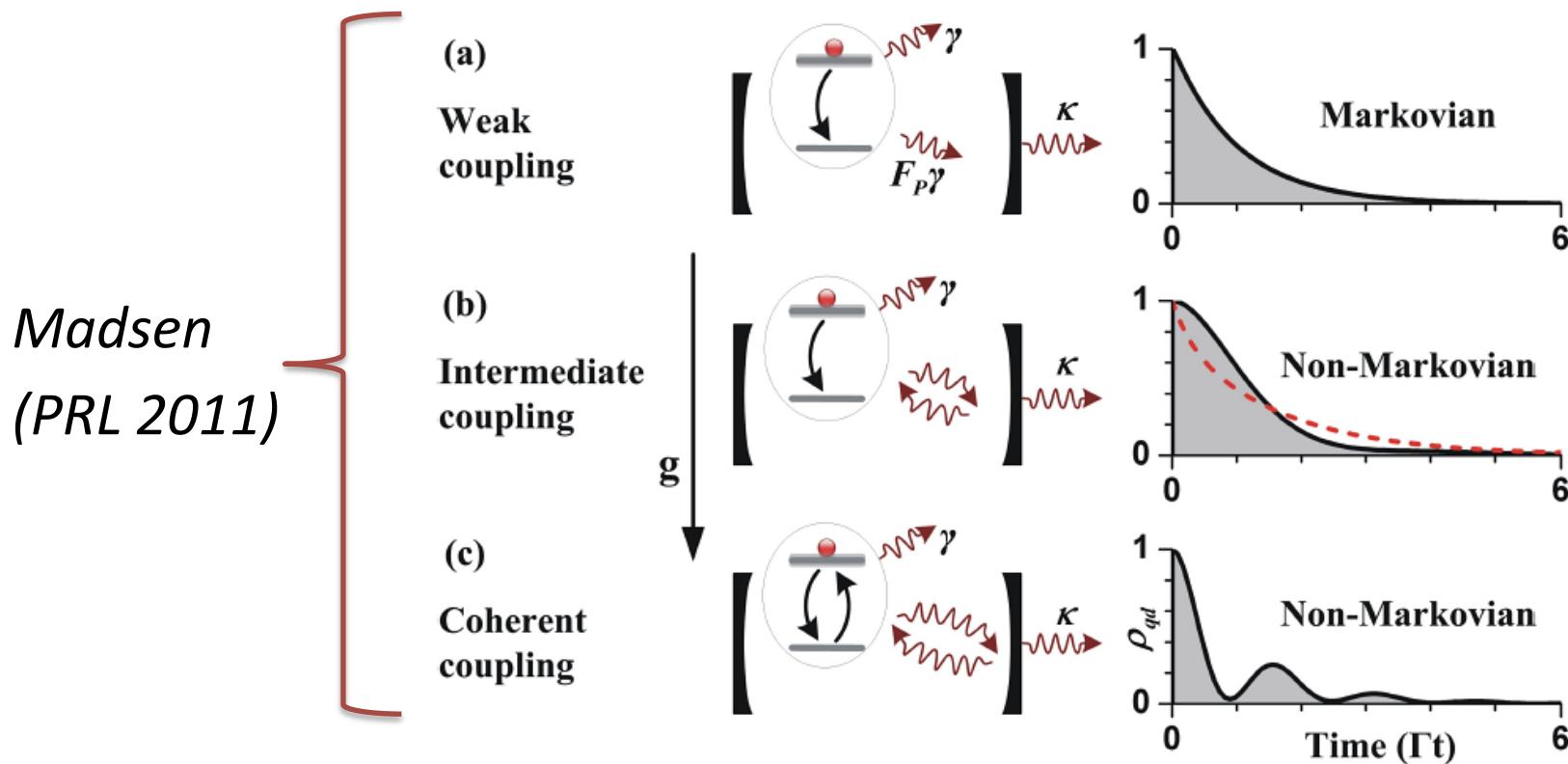
Lian-Ao
(PRE 2009)

$$J = \frac{1}{2} \sum_{n,m} E_{m,n} |S_{m,n}|^2 P_n \times [k_{n \rightarrow m}^L(T_L) - k_{n \rightarrow m}^R(T_R)]$$

- Estudiar la transferencia de calor
- Resultados preliminares

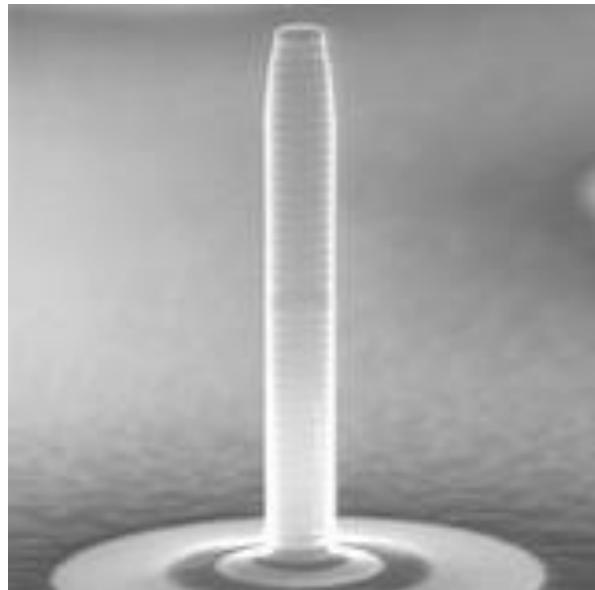
P4: Excitones en un punto cuántico

Caracterizar la dinámica en un sistema de excitones confinados en un punto cuántico



P4: Excitones en un punto cuántico

Caracterizar la dinámica en un sistema de excitones confinados en un punto cuántico



- Cantidadas dinámicas relacionadas con propiedades ópticas medibles
- No existe modelo teórico

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(LPTMC, Paris) L. Cugliandolo (LPTHE, Paris)*