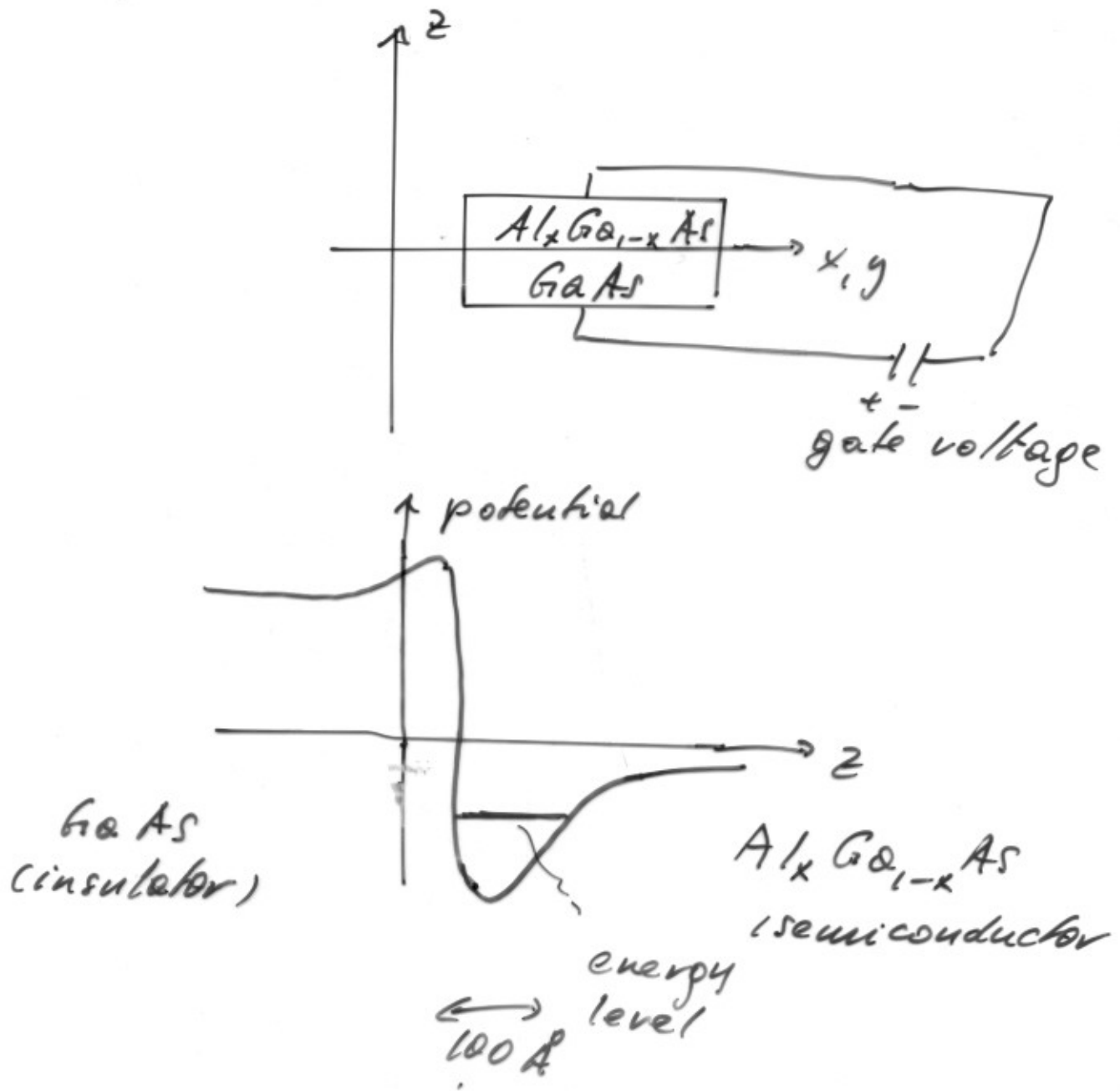


Plan

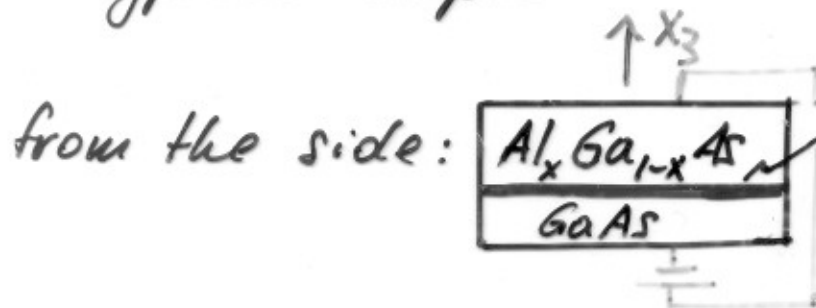
1. The classical Hall effect
2. The quantum Hall effect (QHE)
3. Quantum mechanics: bare essentials
4. Landau levels
5. Stability of the Hall conductance (role of disorder)
6. The topological approach to the QHE
7. The QHE as a pump
8. Edge currents
9. Mathematical reformulations of 6-8
10. Equivalences

The quantum Hall effect : experimental setup

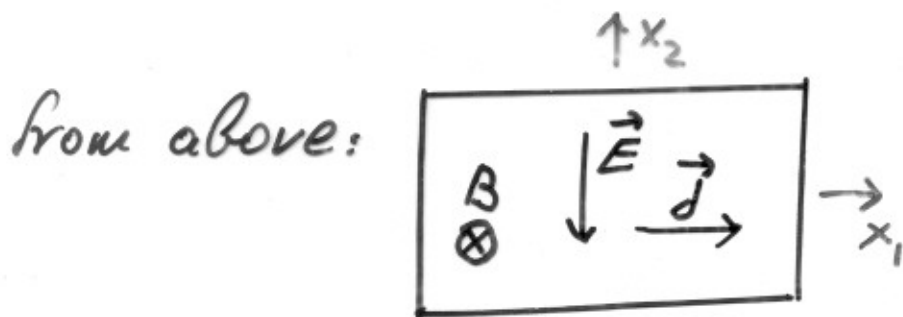


The experiment (von Klitzing, 1980)

a typical sample



- electron gas confined to the interface (dim = 2)
- of density n (or Fermi energy μ) tunable through gate voltage

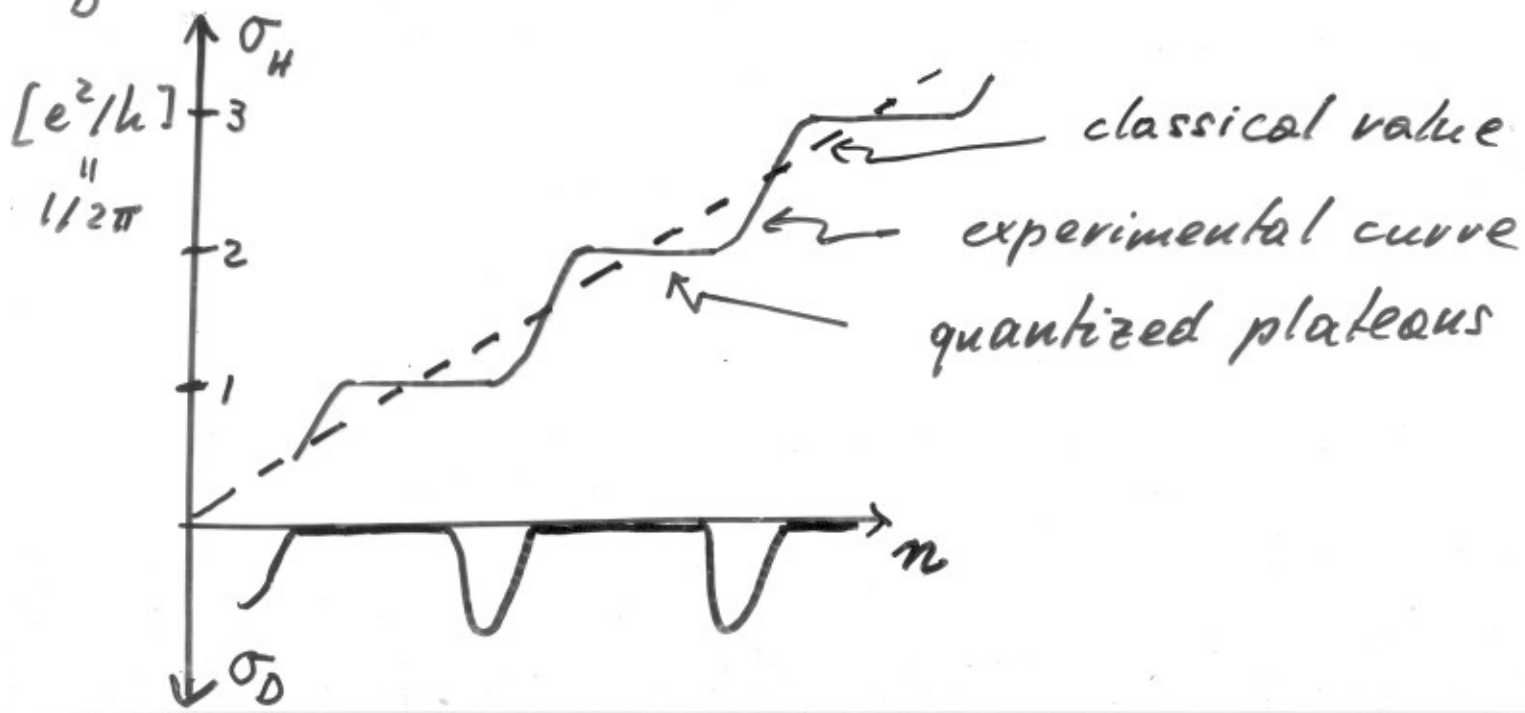


Hall-Ohm law

$$\vec{j} = \sigma \vec{E}, \quad \sigma = \begin{pmatrix} \sigma_D & -\sigma_H \\ \sigma_H & \sigma_D \end{pmatrix}$$

σ_H : Hall conductance

σ_D : ohmic (dissipative) conductance



The bare essentials of quantum mechanics

A quantum mechanical system is characterized by

- a Hilbert space \mathcal{H} (with inner product (\cdot, \cdot))
- some self-adjoint operators A
- a distinguished self-adjoint operator H (the Hamiltonian), possibly time-dependent.

meaning:

- $\psi \in \mathcal{H}$ with $(\psi, \psi) = 1$ is the state of the system at a given time (with ψ and $e^{i\alpha} \psi$, $(\alpha \in \mathbb{R})$ representing the same state)

- A are observables:

$$(\psi, A\psi)$$

is the expectation value of the measurement of A in the state ψ

- Energy H generates evolution ($H = H(t)$ if system subject to external forces) through

$$i\hbar \frac{\partial \psi}{\partial t} = H(t) \psi(t);$$

equivalently: propagator $U(t, t') : \psi(t') \mapsto \psi(t)$

$$i\hbar \frac{\partial}{\partial t} U(t, t') = H(t) U(t, t'), \quad U(t, t) = 1$$

Example: Single-particle systems

- $\mathcal{H} = L^2(\mathbb{R}^d)$, $d = \text{dimension of physical space}$

states ψ "wave functions"

- Observables x_i : position ($i=1, \dots, d$)

$$p_i = -i\hbar \frac{\partial}{\partial x_i} \quad \text{momentum}$$

and, for $d=3$, $\vec{L} = \vec{x} \wedge \vec{p}$ angular momentum

.....

- 1) $H = \frac{\vec{p}^2}{2m} + V(\vec{x}) = -\frac{\hbar^2}{2m} \Delta + V(\vec{x})$

particle in a potential V (e.g. harmonic oscillator, hydrogen atom)

- 2) $A_0(\vec{x}, t), \vec{A}(\vec{x}, t)$ electromagnetic potentials

$$\vec{E} = -\vec{\nabla} A_0 - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \quad \text{electric field}$$

$$\vec{B} = \text{rot } \vec{A} \quad (d=3), \quad B = \frac{\partial A_2}{\partial x_1} - \frac{\partial A_1}{\partial x_2} \quad (d=2)$$

magnetic field

$$\frac{1}{m} \left(\vec{p} - \frac{e}{c} \vec{A}(\vec{x}, t) \right) : \text{velocity}$$

$$H(t) = \frac{1}{2m} \left(\vec{p} - \frac{e}{c} \vec{A}(\vec{x}, t) \right)^2 + A_0(\vec{x}, t)$$

Example: Systems of many non-interacting fermionic particles

may be described in terms of corresponding single-particle system \mathcal{H} .

- Pauli principle: Each single-particle state can be occupied at most once

- Projection P onto the occupied state $\psi_i \in \mathcal{H}$ defines a many-particle state

- Expectation of A in $P = \sum_i \psi_i (\psi_i, \cdot)$

$$\text{tr}(AP) = \sum_i (\psi_i, A \psi_i)$$

- Evolution from t' to t

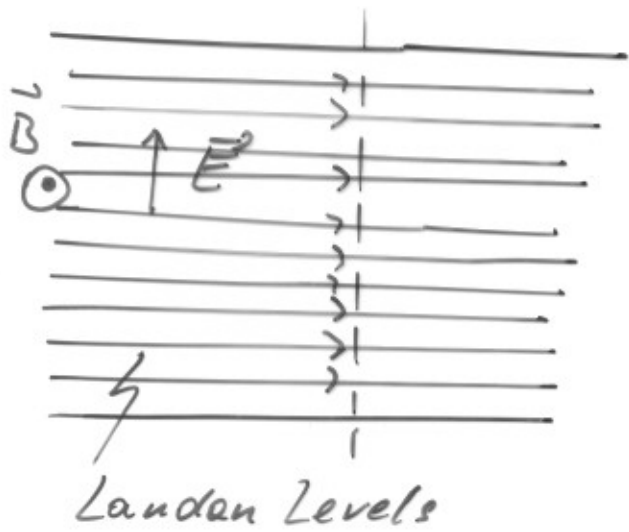
$$P \mapsto U(t, t') P U(t, t')^*$$

- Many-body ground state

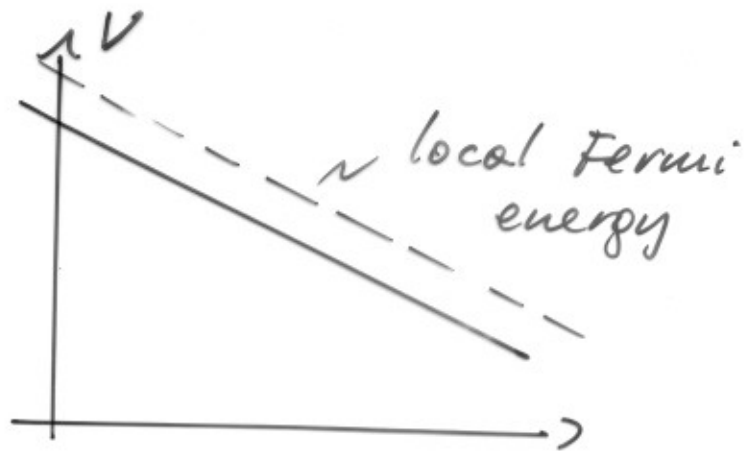
P is the projection onto the eigenstates of H with lowest eigenvalues (energies) up to some energy μ (Fermi energy)

μ $\sigma(H)$: spectrum of H

- pure sample



cross section



- disordered sample

